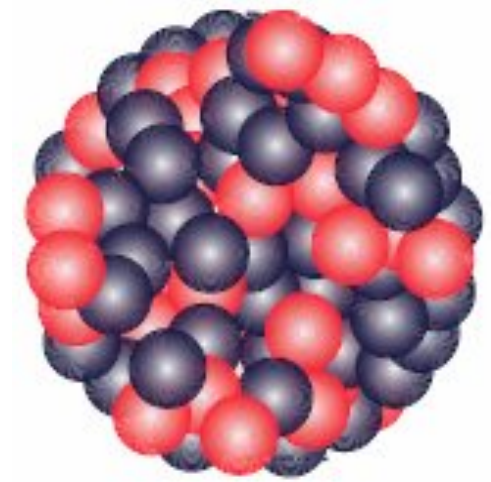


8 - *Nuclei*

8.1 - The nucleus

The atomic nucleus consists of **protons** and **neutrons**.

Protons and neutrons are called **nucleons**.



A nucleus is characterized by:

- **A**: Mass Number = number of nucleons
- **Z**: Charge Number = number of protons
- **N**: Neutron Number

Determines the element

Determines the isotope

Of course $A=Z+N$

Usual notation:

Mass number **A**

^{12}C

Element symbol - defined by charge number
C is Carbon and **Z** = 6

So this nucleus is made of 6 protons and 6 neutrons

8.1.1 - Nuclear physics

a. Nucleons size: ~1 fm

b. Nuclei

nucleons attract each other via the strong force (range ~ 1 fm)

a bunch of nucleons bound together create a potential for an additional :

	Mass	Spin	Charge
Proton	938.272 MeV/c ²	1/2	+e
Neutron	939.565 MeV/c ²	1/2	0

neutron

proton

(or any other charged particle)

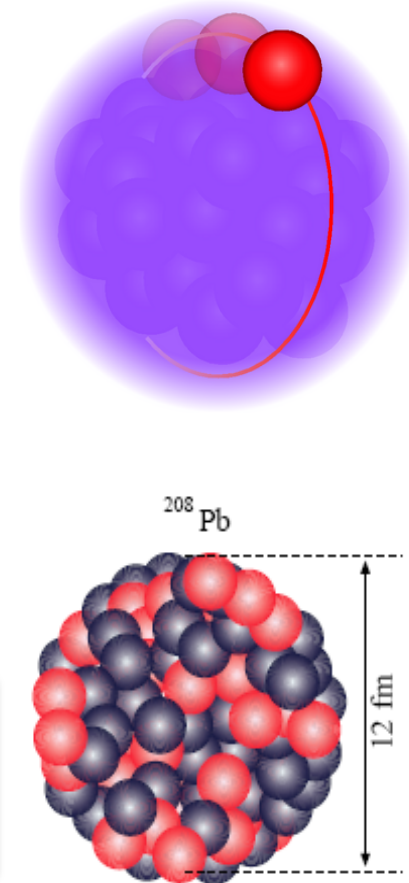
$$R \sim 1.3 \times A^{1/3} \text{ fm}$$

Coulomb Barrier V_c

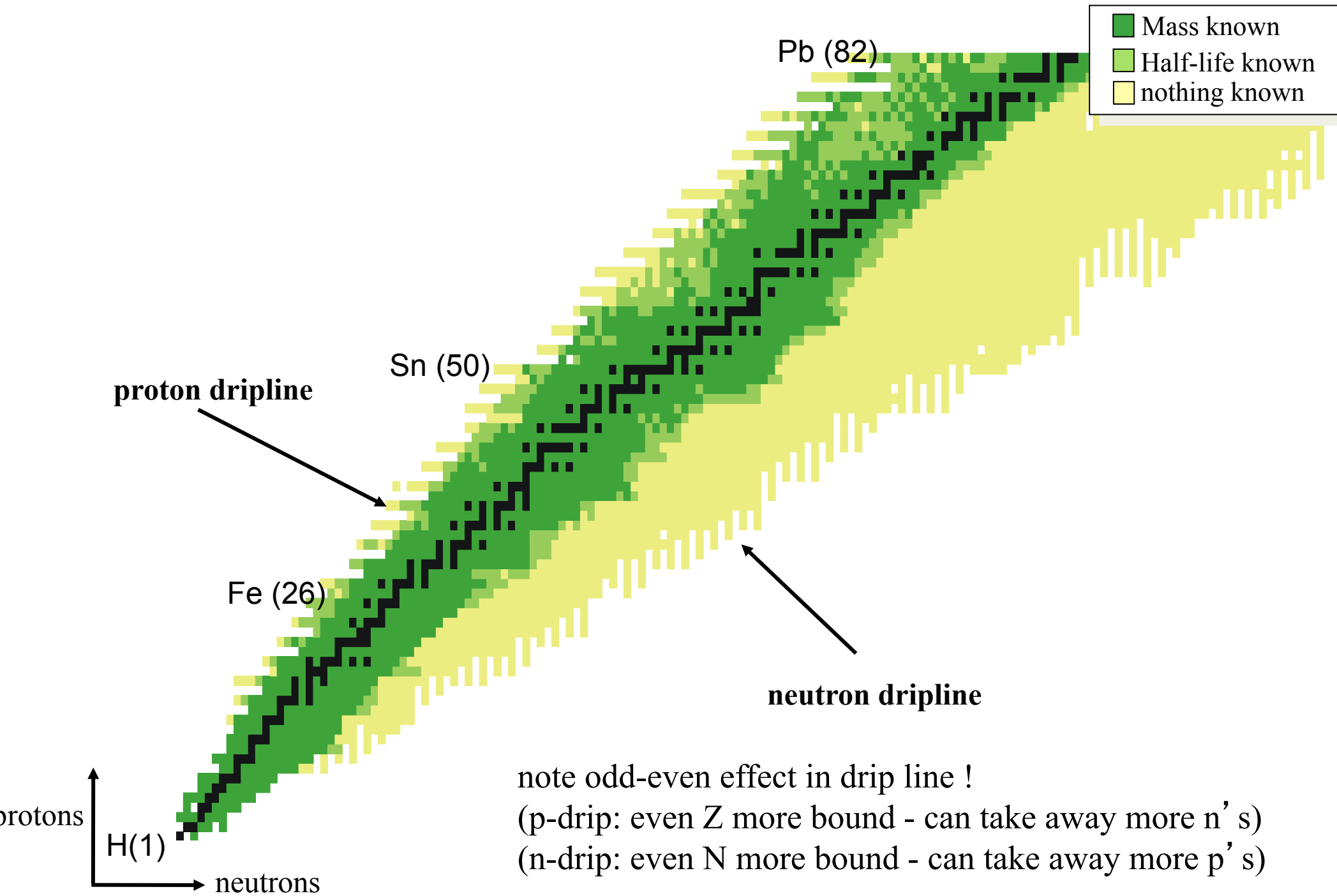
Nucleons in a Box:
Discrete energy levels
in nucleus

$$V_c = \frac{Z_1 Z_2 e^2}{R}$$

(8.1)



The Nuclear Chart



8.1.2 - Nuclear Masses and Binding Energies

Energy that is released when a nucleus is assembled from neutrons and protons

$$m(Z, N) = Zm_p + Nm_n - B / c^2 \quad (8.2)$$

m_p = proton mass, m_n = neutron mass, $m(Z, N)$ = mass of nucleus with Z, N

- $B > 0$
- With B the mass of the nucleus is determined.
- B is roughly $\sim A$

Masses are usually tabulated as atomic masses

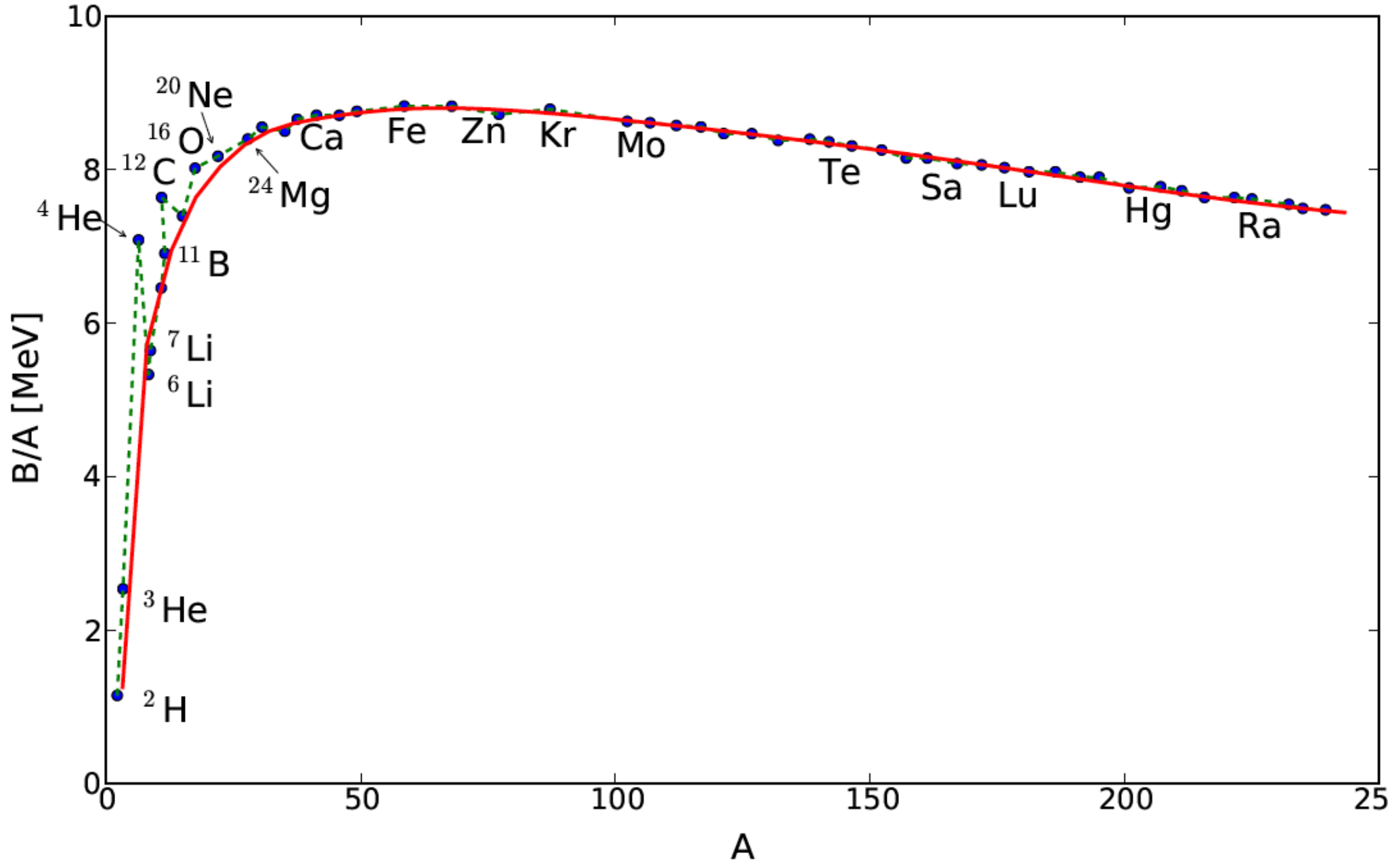
$$m = m_{nuc} + Z m_e + B_e \quad (8.3)$$



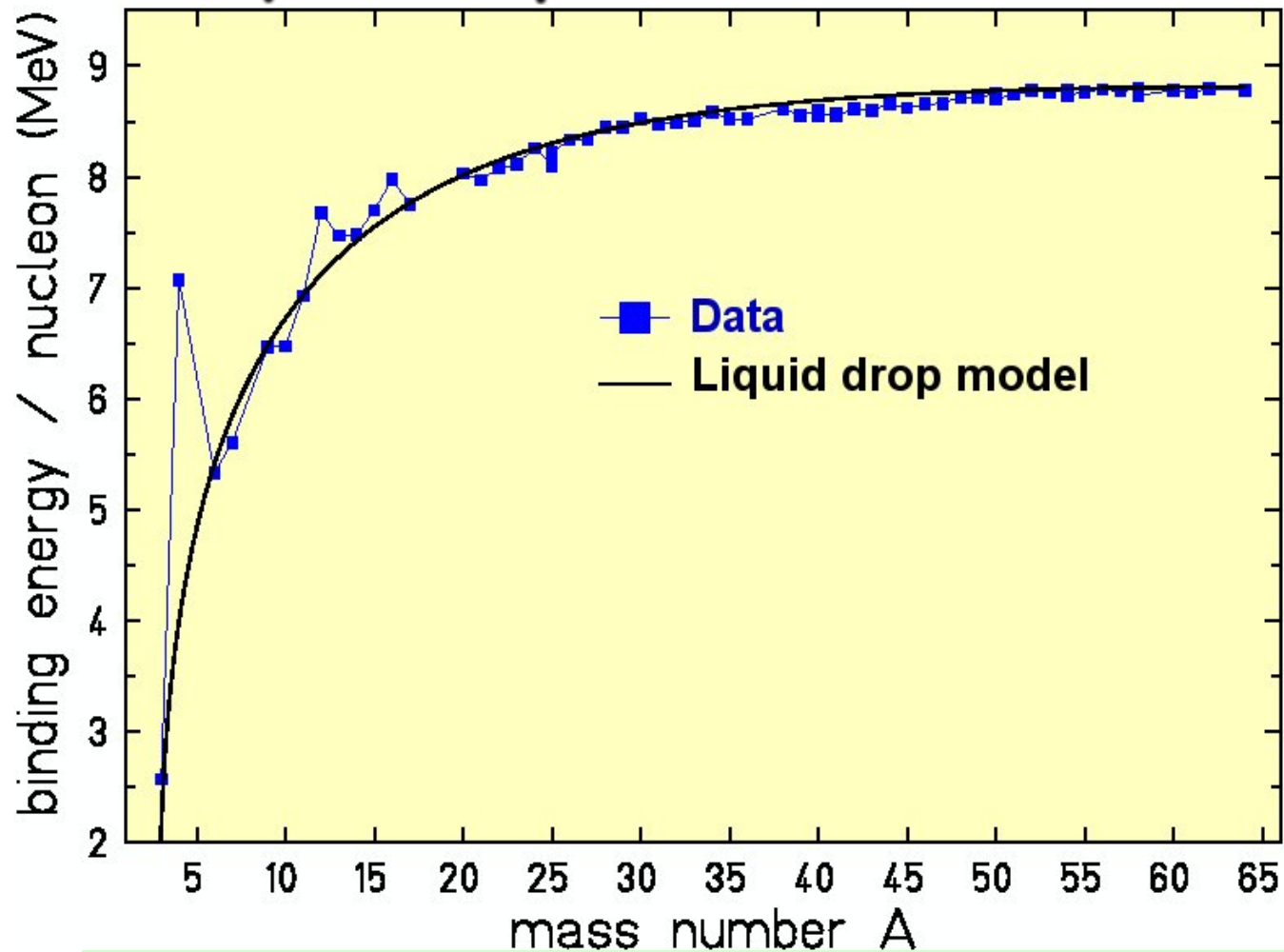
Most tables give atomic mass excess Δ in MeV:
 (definition: for ^{12}C : $\Delta = 0$)

$$m = Am_u + \Delta / c^2 \quad (8.4)$$

Nuclear Masses and Binding Energies



The Liquid Drop Model



Bethe-Weizsäcker formula

$$B(Z, N, A) = a_V A - a_{Surf} A^{2/3} - a_{sym} \frac{(Z - N)^2}{A} + a_{Coul} Z(Z - 1) A^{-1/3} \quad (8.5)$$

$$a_V = 15.85 \text{ MeV}, \quad a_{Surf} = 18.34 \text{ MeV}, \quad a_{Symm} = 23.21 \text{ MeV}, \quad a_{Coul} = 0.71 \text{ MeV} \quad (8.6)$$

Understanding the Liquid Drop Model

Assumes incompressible fluid (volume $\sim A$) and sharp surface)

$$B(Z, A) = a_V A \quad \text{Volume Term} \quad (8.7)$$

$$(8.8) \quad -a_s A^{2/3} \quad \text{Surface Term} \quad \sim \text{surface area (Surface nucleons less bound)}$$

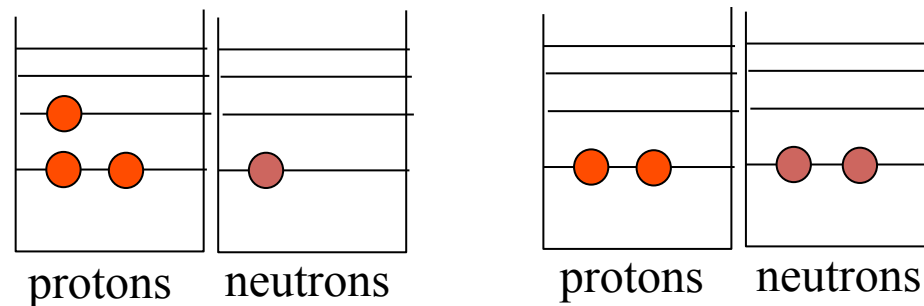
$$(8.9) \quad -a_C \frac{Z^2}{A^{1/3}} \quad \text{Coulomb term. Coulomb repulsion leads to reduction uniformly charged sphere has}$$

$$E = \frac{3}{5} \frac{Q^2}{R}$$

$$-a_A \frac{(Z - A/2)^2}{A}$$

Asymmetry term: Pauli principle to protons: symmetric filling of p,n potential boxes has lowest energy (ignore Coulomb)

(8.10)



lower total energy = more bound

(8.11)

$$+a_p A^{-1/2} \quad \left\{ \begin{array}{ll} \times 1 & ee \\ \times 0 & oe, eo \\ \times (-1) & oo \end{array} \right. \quad \text{Pairing term: even number of like nucleons favored}$$

(e=even, o=odd referring to Z, N respectively)

Understanding the Liquid Drop Model

neglect asymmetry term (assume reasonable asymmetry)

neglect pairing and shell corrections (see later) - want to understand average behavior

then

$$B / A = a_V - a_S \frac{1}{A^{1/3}} - a_C \frac{Z^2}{A^{4/3}} \quad (8.12)$$

const
as strong force
has short range

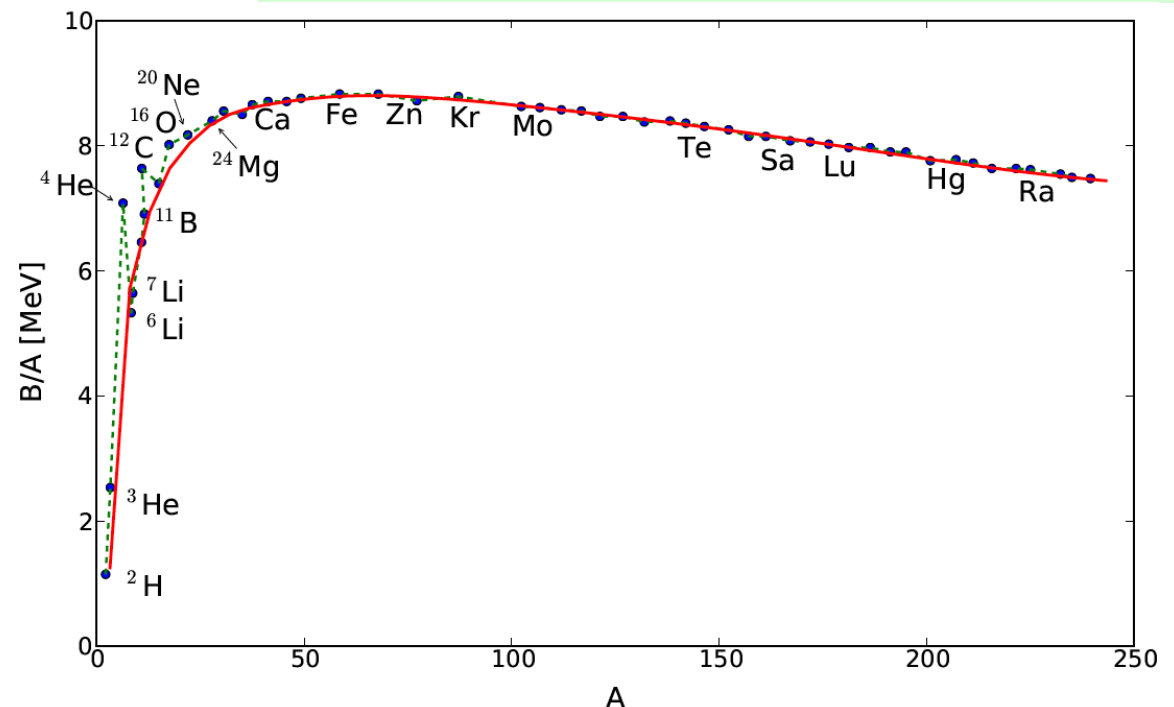
~surface/volume ratio
favors large nuclei

Coulomb repulsion has long range
- the more protons the more repulsion
favors small (low Z) nuclei

Maximum around ~ Fe

→ **Fusion** of light elements
release nuclear energy

→ **Fission** of very heavy elements
also release nuclear energy

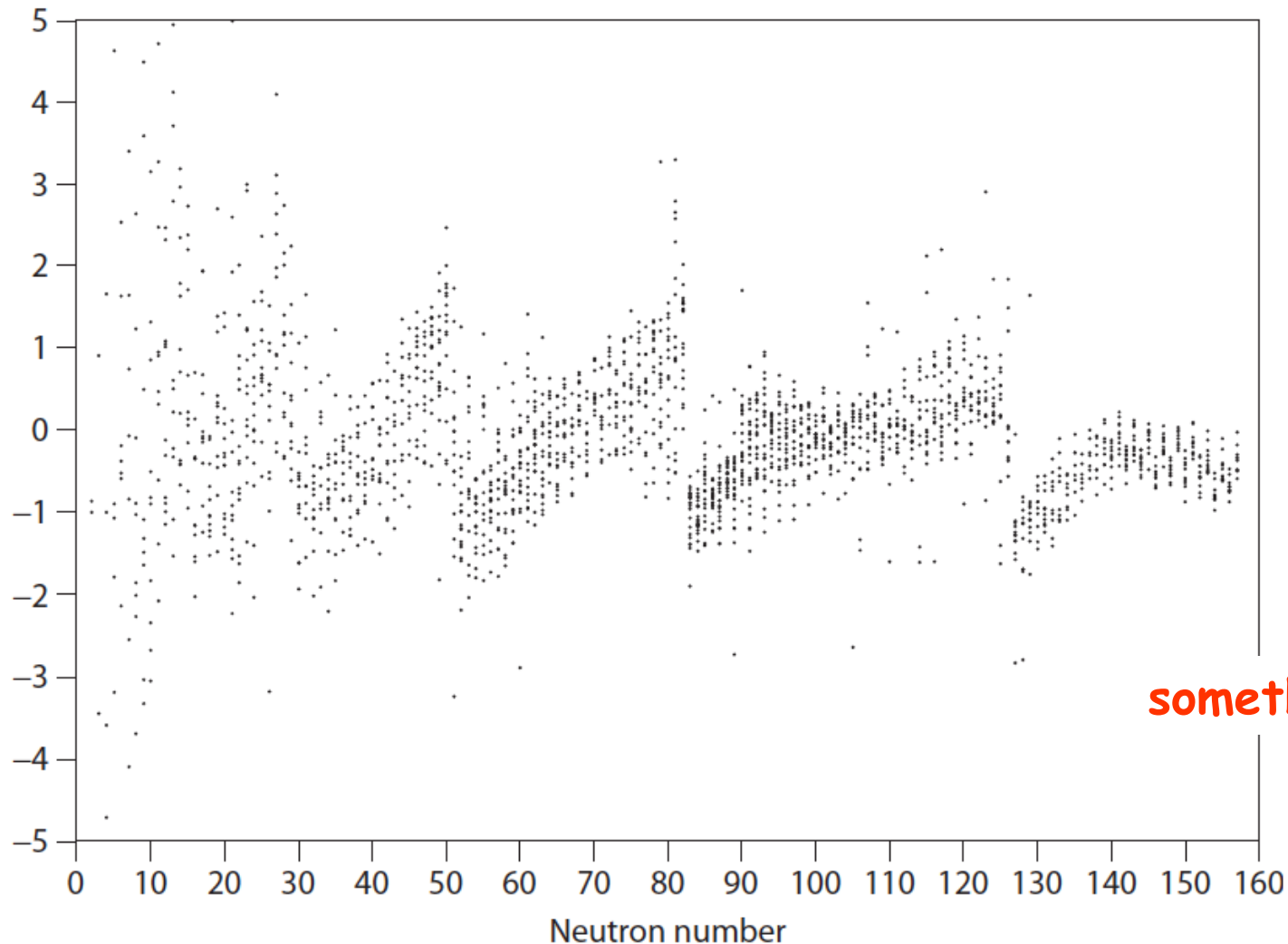


Fitting the Liquid Drop Model to Experimental Data

in MeV/c²

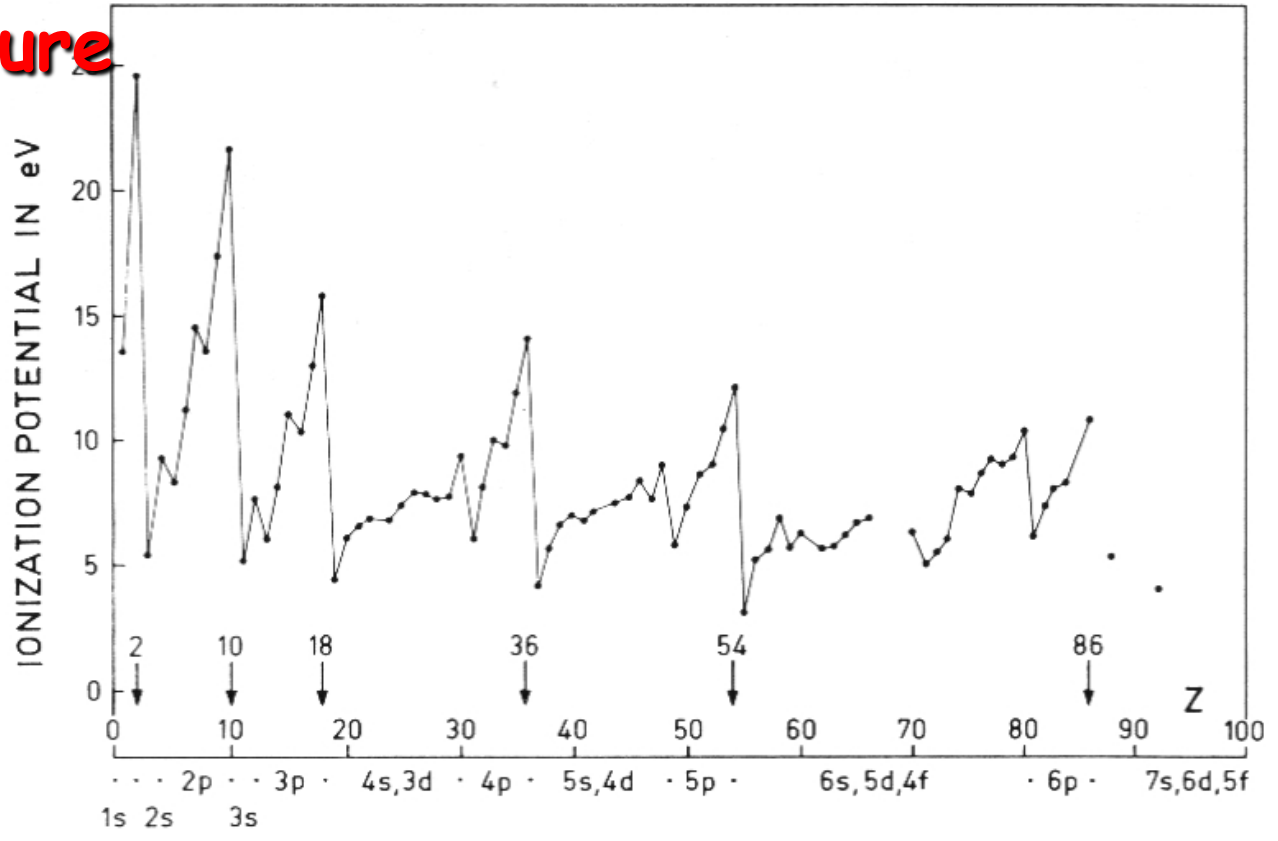
a_V	a_S	a_C	a_A	a_P
15.85	18.34	0.71	92.86	11.46

Deviation (in MeV) to experimental masses:

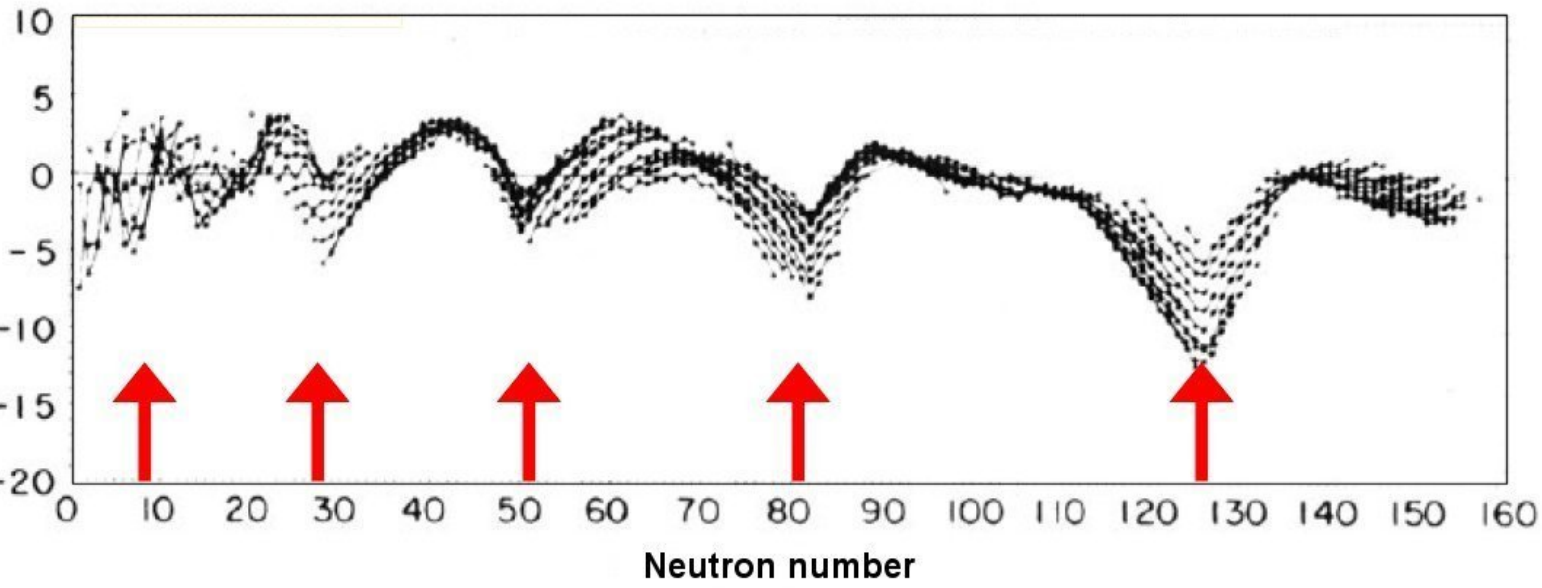


8.1.2 - Shell Structure

Atomic physics



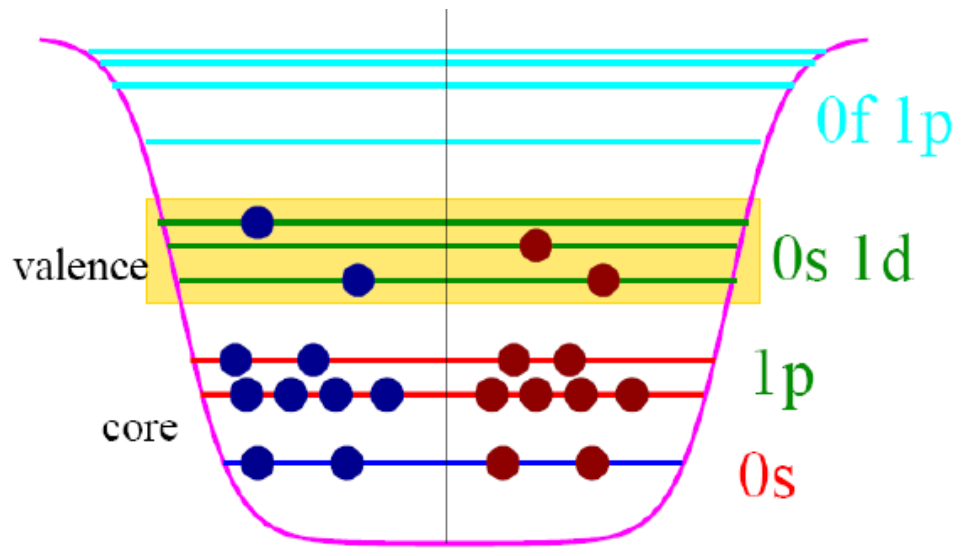
$B_{exp} - B_{LDM}$



Nuclear physics

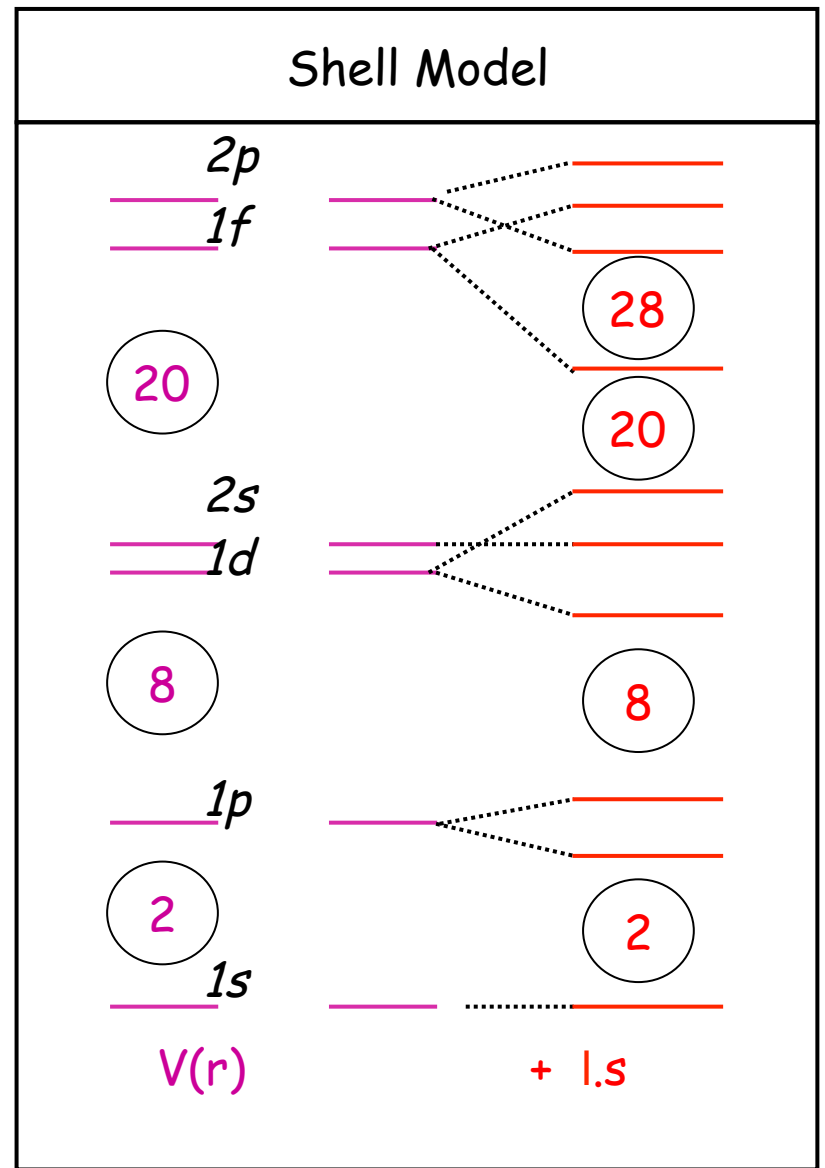
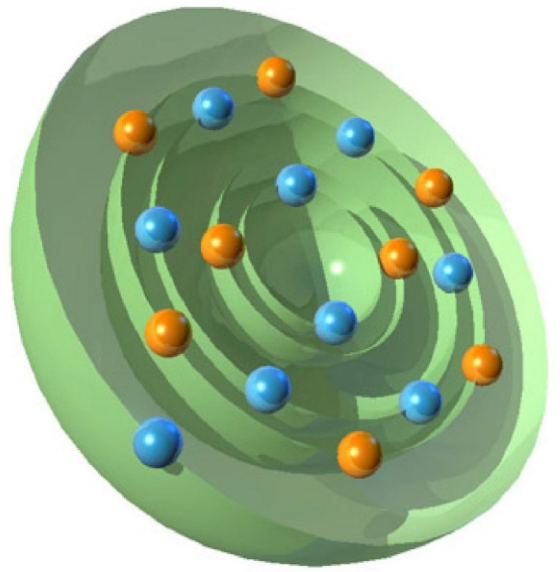
$N_{magic} = 2, 8, 20, 28, 50, 82, 126$

Shell Structure Model



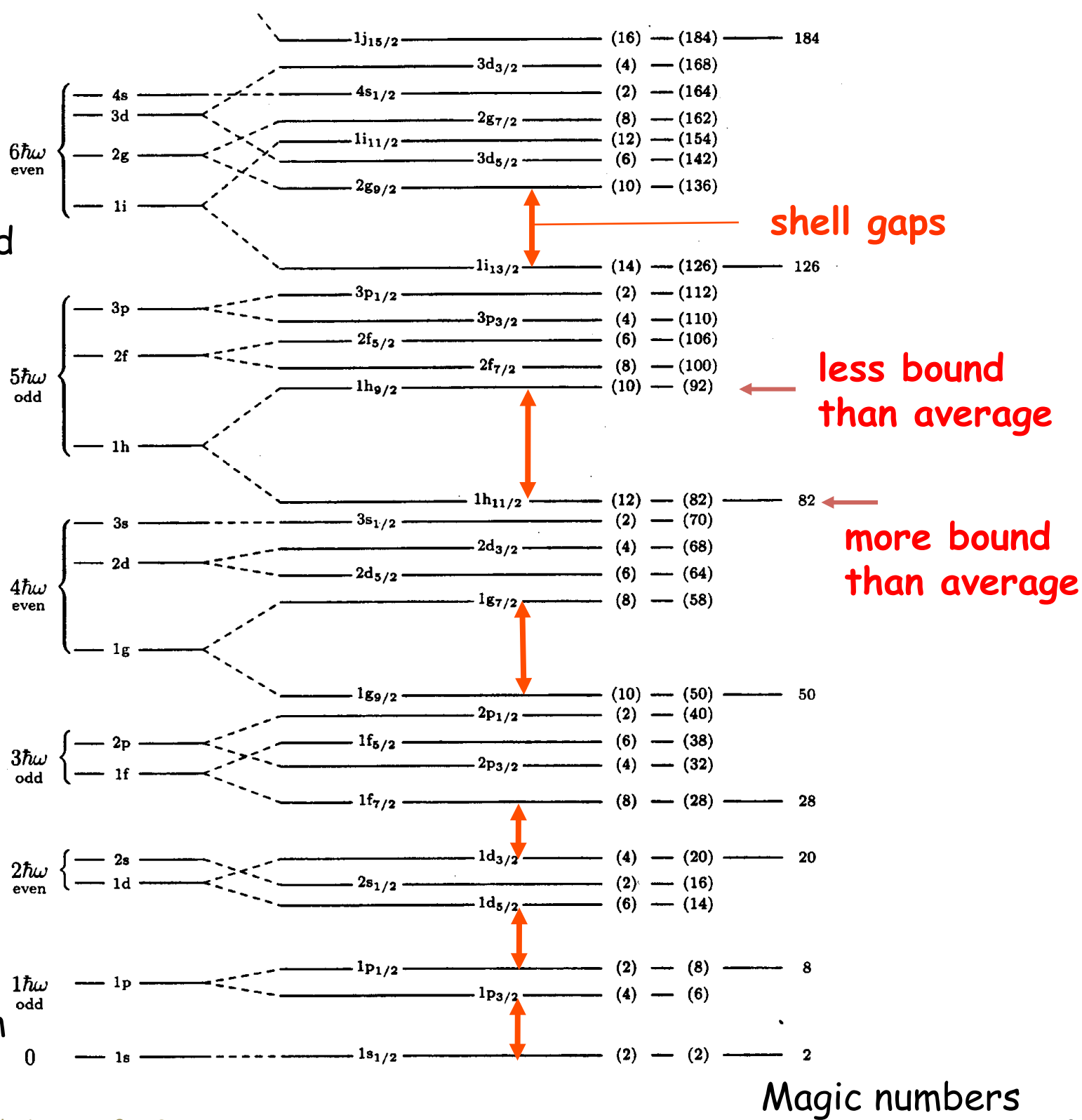
$$H = \sum_{k=1}^A \left[\frac{p^2}{2m} + V(r) + V_{SO}(r) \mathbf{l} \cdot \mathbf{s} \right] \quad (8.13)$$

Mayer, Jensen, Nobel Prize, 1963



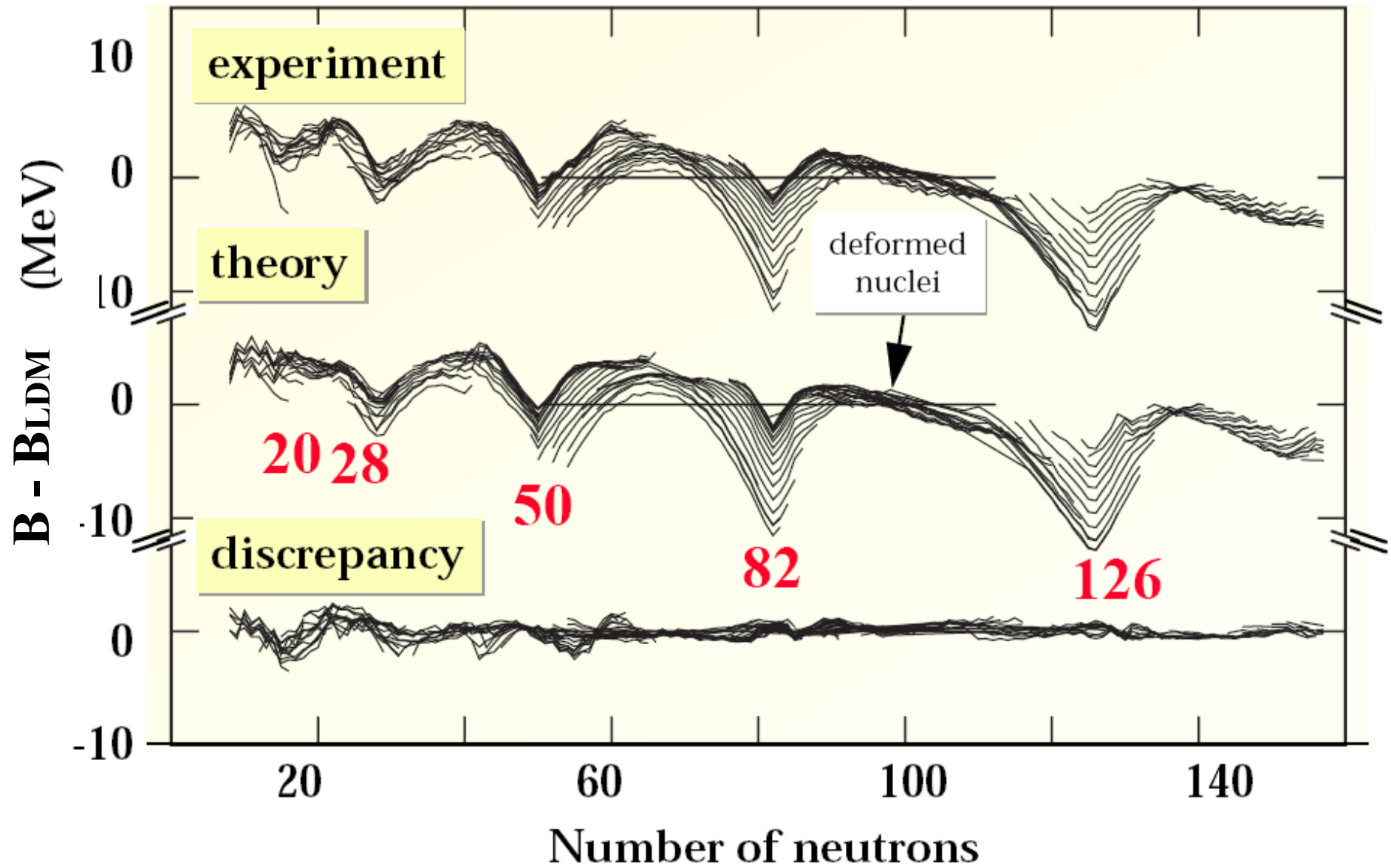
Shell model: (single nucleon energy levels)

are not evenly spaced



need to add
shell correction term
 $S(Z,N)$

Liquid Drop Formula corrected with shell model



8.2 - Nuclear decay

a) Energy generation

nuclear reaction: $A + B \longrightarrow C$

if $m_A + m_B > m_C$ then energy $Q = (m_A + m_B - m_C)c^2$ is generated by reaction
(7.14)

“Q-value” Q = Energy generated (>0) or consumed (<0) by reaction

b) Stability

if there is a reaction $A \longrightarrow B + C$ with $Q > 0$ (or $m_A > m_B + m_C$)
then decay of nucleus A is energetically possible.

nucleus A might then not exist (at least not for a very long time)

Nuclear decay

Decay of A in B and C is possible if reaction $A \longrightarrow B + C$ has positive Q -value

BUT: there might be a barrier that prolongs the lifetime

Decay is described by quantum mechanics and is a pure random process, with a constant probability for the decay to happen in a given time interval.

N : Number of nuclei A (parent)

λ : decay rate (decays per second and parent nucleus)

$$dN = -\lambda N dt$$

(8.15)

therefore

$$N(t) = N(t = 0) e^{-\lambda t}$$

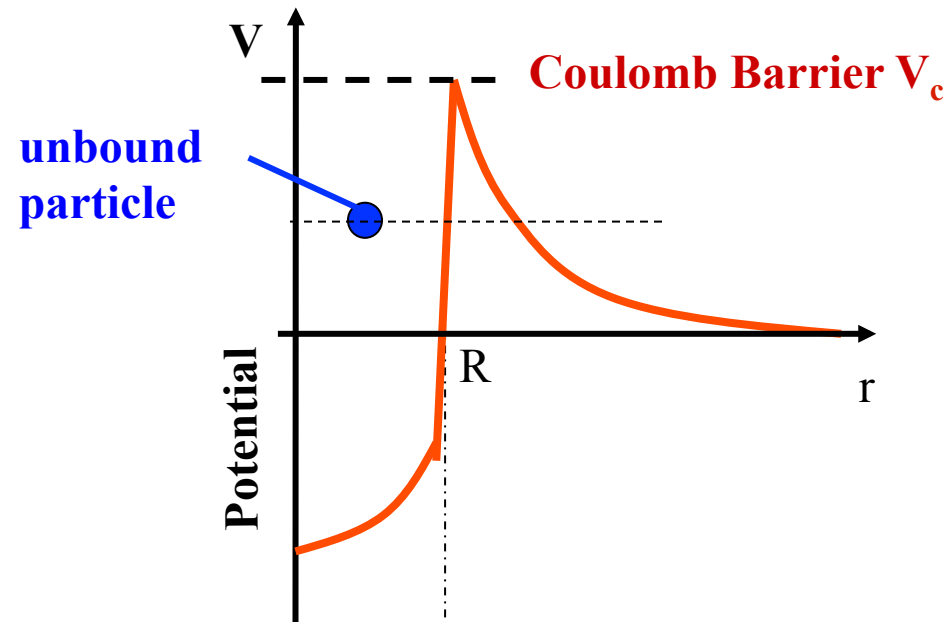
(8.16)

lifetime $\tau = 1/\lambda$

half-life $T_{1/2} = \tau \ln 2 = \ln 2 / \lambda$ is time for half of the nuclei present to decay

Decay modes

for anything other than a neutron (or a neutrino) emitted from the nucleus there is a Coulomb barrier



$$V_c = \frac{Z_1 Z_2 e^2}{R}$$

If that barrier delays the decay beyond the lifetime of the universe (~ 14 Gyr) we consider the nucleus as being stable.

Example: for $^{197}\text{Au} \rightarrow ^{58}\text{Fe} + ^{139}\text{I}$ has $Q \sim 100 \text{ MeV}$!
yet, gold is stable.



not all decays that are energetically possible happen

most common:

- β decay
- n decay
- p decay
- α decay
- fission

8.2.1 - β decay

$p \leftrightarrow n$ conversion within a nucleus via weak interaction

Modes (for a proton/neutron in a nucleus):

$$\beta^+ \text{ decay} \quad p \longrightarrow n + e^+ + \nu_e$$

$$\text{electron capture} \quad e^- + p \longrightarrow n + \nu_e$$

$$\beta^- \text{ decay} \quad n \longrightarrow p + e^- + \bar{\nu}_e$$

(8.17)

Electron capture (or EC) of atomic electrons or, in astrophysics, of electrons in the surrounding plasma

Q-values for decay of nucleus (Z,N)

with nuclear masses

with atomic masses

$$Q_{\beta^+} / c^2 = m_{\text{nuc}}(Z,N) - m_{\text{nuc}}(Z-1,N+1) - m_e = \mathbf{m(Z,N) - m(Z-1,N+1) - 2m_e}$$

$$Q_{\text{EC}} / c^2 = m_{\text{nuc}}(Z,N) - m_{\text{nuc}}(Z-1,N+1) + m_e = \mathbf{m(Z,N) - m(Z-1,N+1)} \quad (8.18)$$

$$Q_{\beta^-} / c^2 = m_{\text{nuc}}(Z,N) - m_{\text{nuc}}(Z+1,N-1) - m_e = \mathbf{m(Z,N) - m(Z+1,N-1)}$$

Note: $Q_{\text{EC}} > Q_{\beta^+}$ by 1.022 MeV

Q-values with Δ values

$$m = Am_u + \Delta / c^2 \quad (8.19)$$

Q-values for reactions that conserve the number of nucleons can also be calculated directly using the tabulated Δ values instead of the masses

Example: $^{14}\text{C} \rightarrow ^{14}\text{N} + e + \nu_e$

$$\begin{aligned} Q / c^2 &= m_{^{14}\text{C}} - m_{^{14}\text{N}} \\ &= 14m_u + \Delta_{^{14}\text{C}} - 14m_u - \Delta_{^{14}\text{N}} \\ &= \Delta_{^{14}\text{C}} - \Delta_{^{14}\text{N}} \quad (\text{for atomic } \Delta' \text{ s}) \end{aligned} \quad (8.20)$$

Q-values with binding energies B

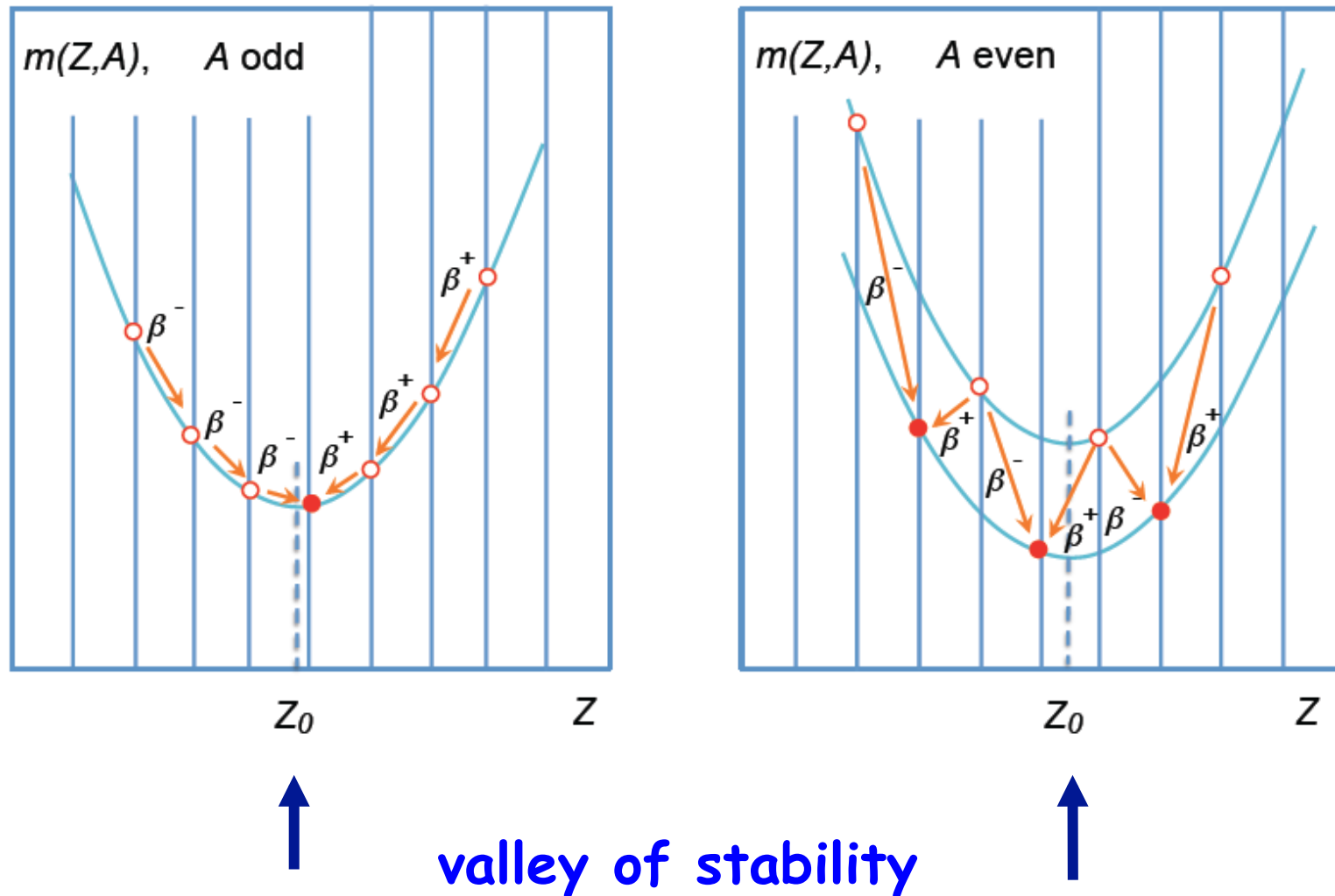
$$m(Z, N) = Zm_p + Nm_n - Bc^2 \quad (8.21)$$

Q-values for reactions that conserve proton number and neutron number can be calculated using $-B$ instead of the masses

β decay

β decay basically no barrier \rightarrow if energetically possible it usually happens (except if another decay mode dominates)

therefore: any nucleus with a given mass number A will be converted into the most stable proton/neutron combination with mass number A by β decays



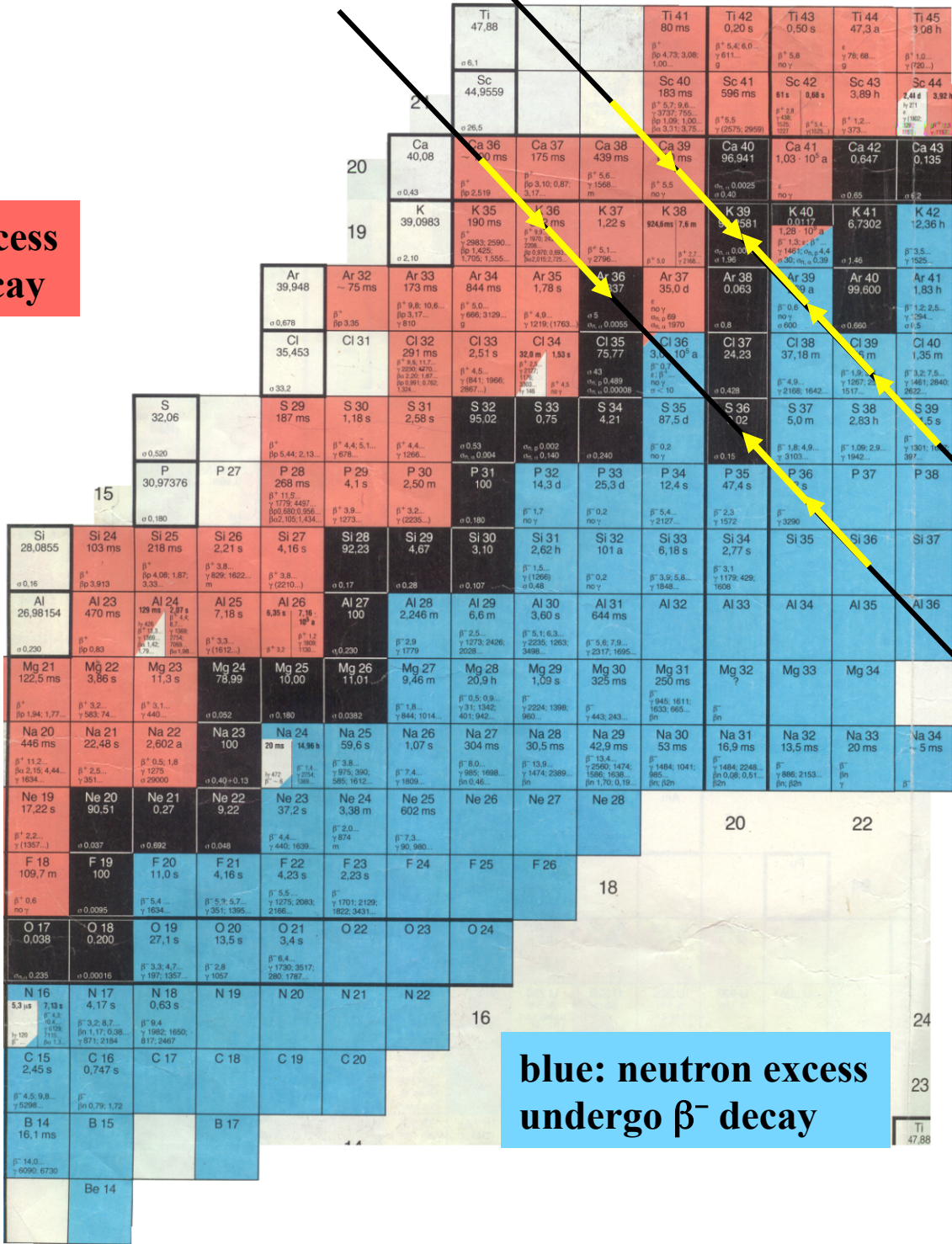
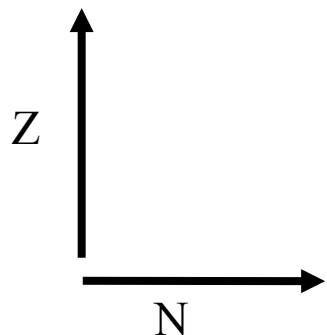
Typical part of the chart of nuclides

red: proton excess
undergo β^+ decay

blue: neutron excess
undergo β^- decay

odd A
isobaric
chain

even A
isobaric
chain



Typical β -decay half-lives

- very near “stability” : occasionally Millions of years or longer
- more common within a few nuclei of stability: minutes - days
- most exotic nuclei that can be formed: ~ milliseconds

